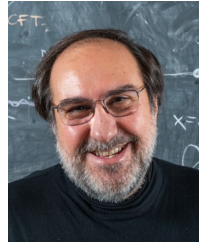


Laurent
Sanchez-Palencia

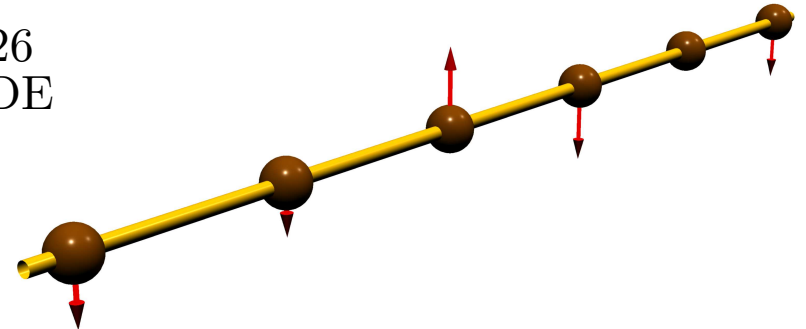


Thierry
Giamarchi

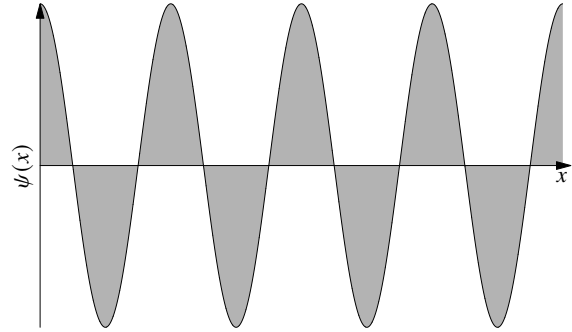
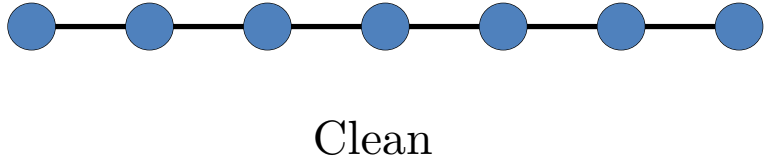
Localization Transition for Interacting Quantum Particles in Colored-Noise Disorder

Giacomo Morpurgo

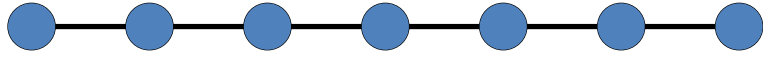
Roscoff, 04/05/2026
5e journées GDR IDE



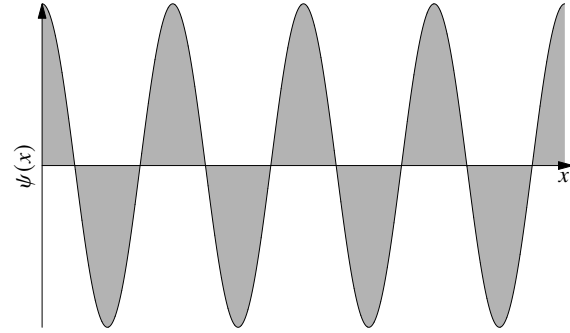
Anderson localization



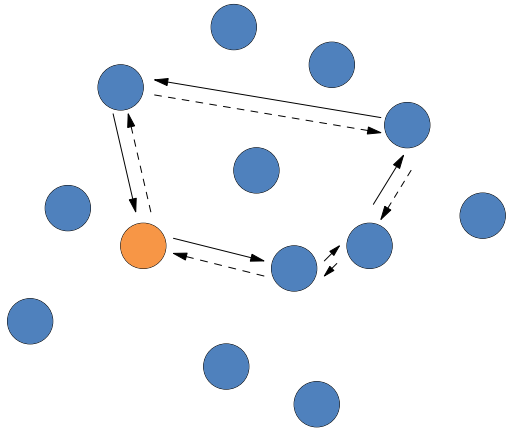
Anderson localization



Clean

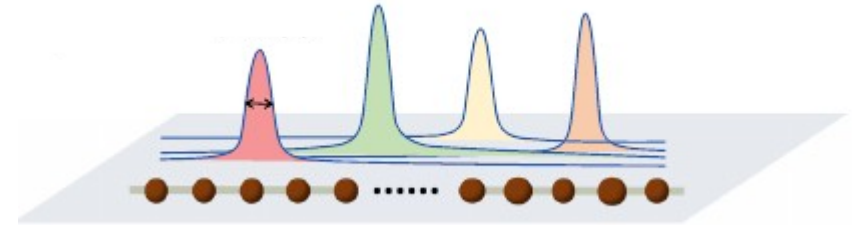
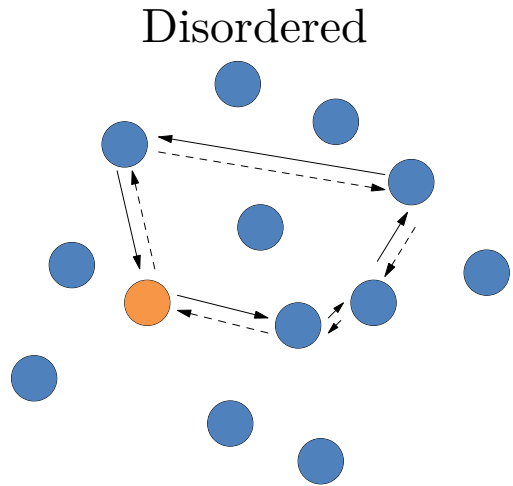
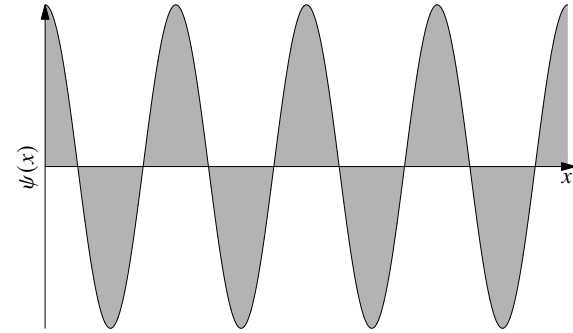
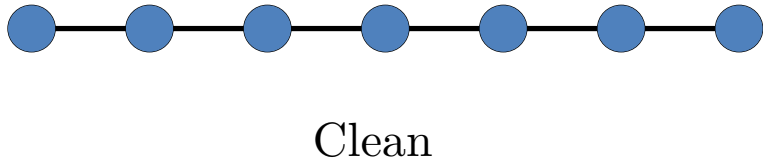


Disordered



- Constructive interference of two opposite paths

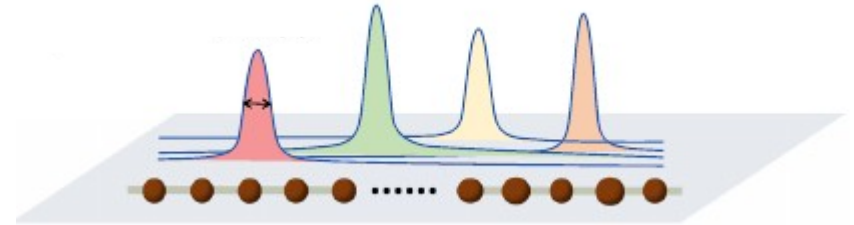
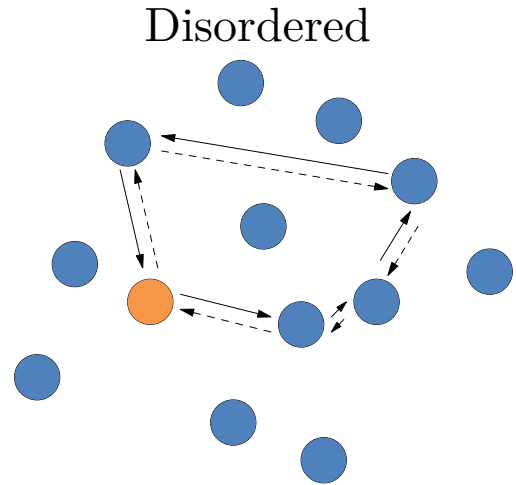
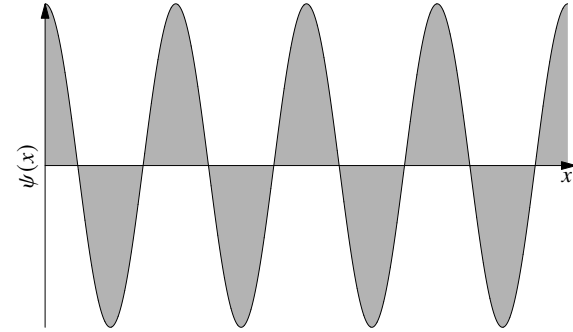
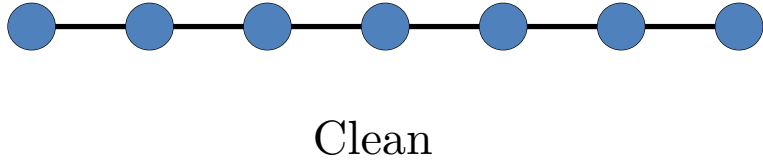
Anderson localization



Guo et al.
Nature Communications 2024

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Anderson localization



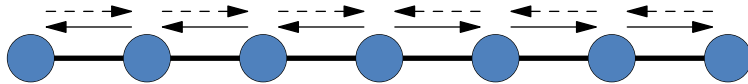
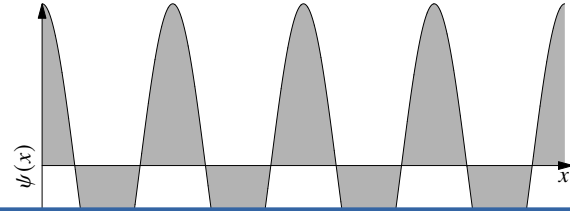
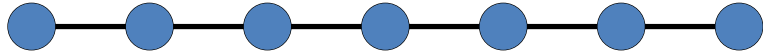
Guo et al.
Nature Communications 2024

Localization length

$$\psi(x) = \frac{1}{\sqrt{\xi}} e^{-\frac{|x-x_0|}{\xi}}$$

- Constructive interference
of two opposite paths

Anderson localization



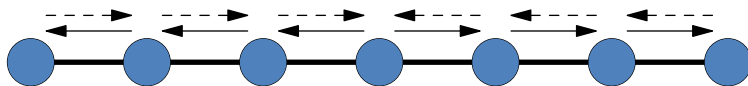
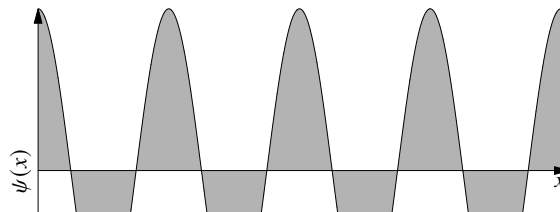
In 1D, infinitesimal disorder
is enough

- Constructive interference
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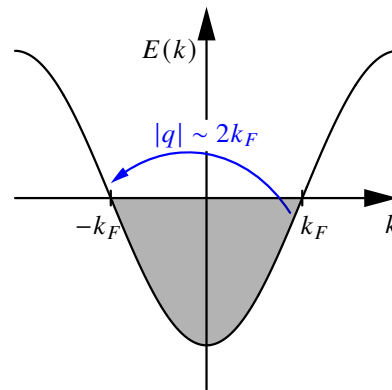
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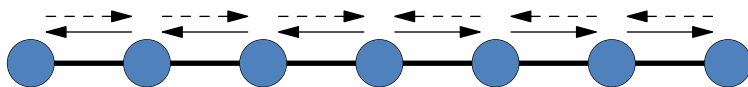
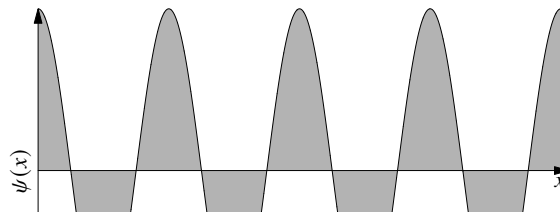


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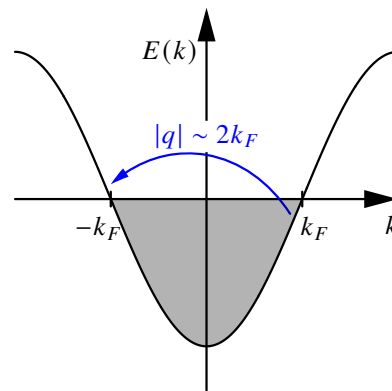
- Constructive interference
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Anderson localization



In 1D, infinitesimal disorder
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$$\xi \propto \frac{1}{W_0^2}$$



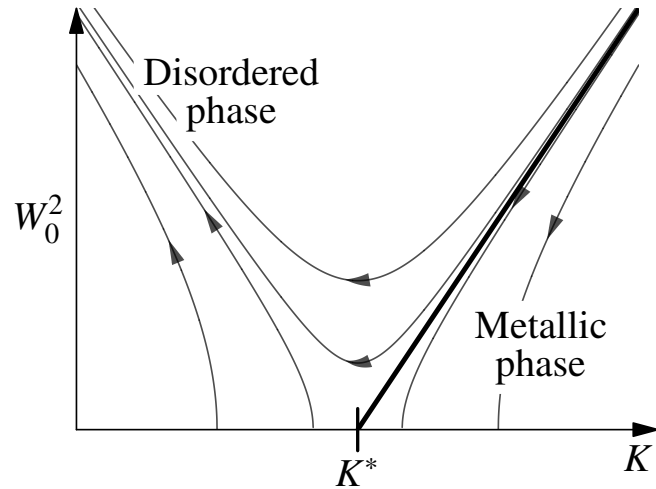
Localization length

$$\psi(x) = \frac{1}{\sqrt{\xi}} e^{-\frac{|x-x_0|}{\xi}}$$

- Constructive interference
of two opposite paths

Adding interactions

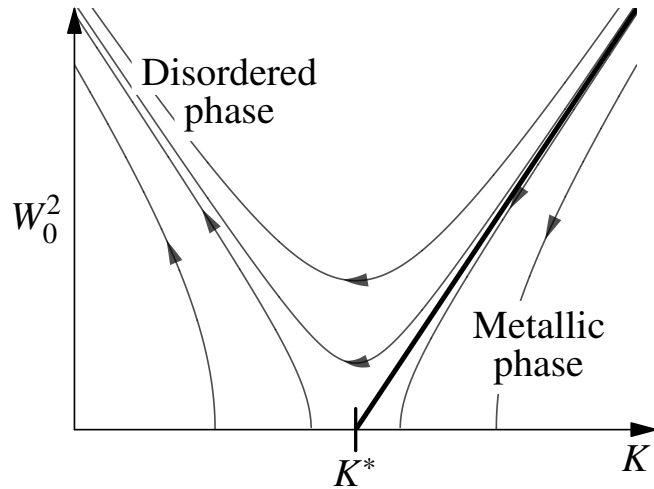
Phase diagram



$$K^* = 3/2$$

Adding interactions

Phase diagram



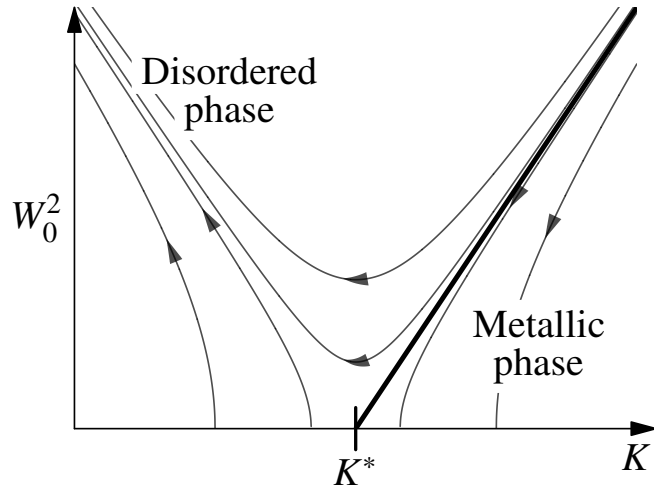
$$K^* = 3/2$$

Localization length

$$\xi \propto \frac{1}{W_0^{2(3-2K)}}$$

Adding interactions

Phase diagram



$$K^* = 3/2$$

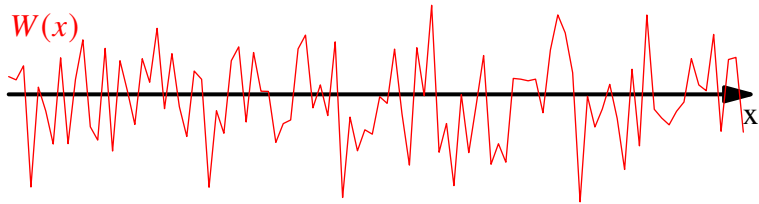
Localization length

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Significant others
-Many body localization
-Bose Glass

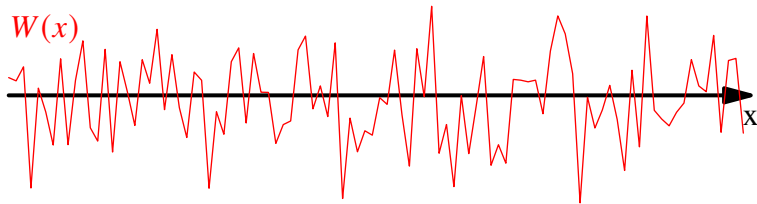
Correlations in the disorder

Uncorrelated

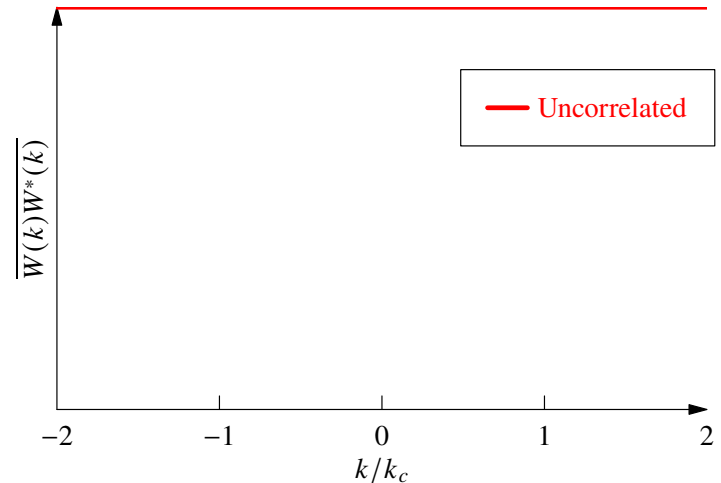
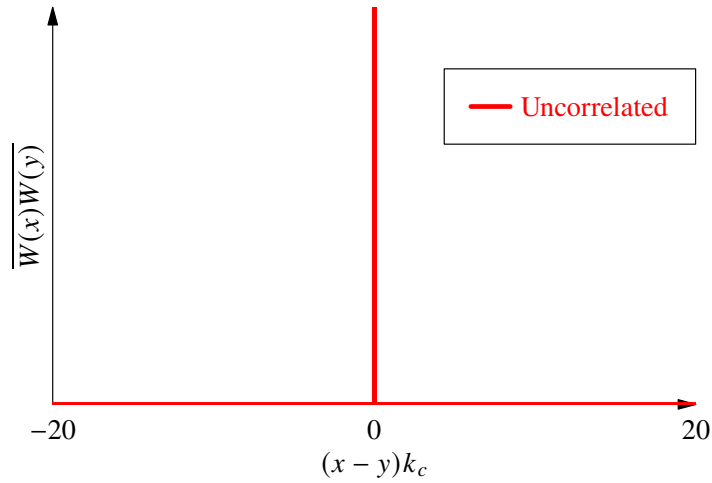


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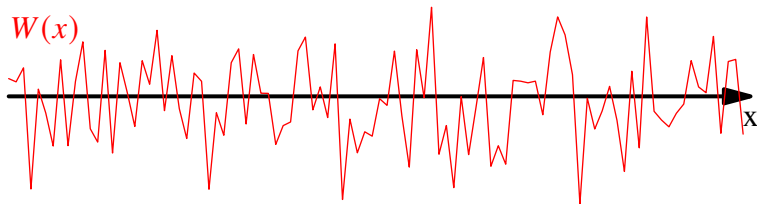


$$\overline{W(k)W^*(k)} = W_0^2\Omega$$



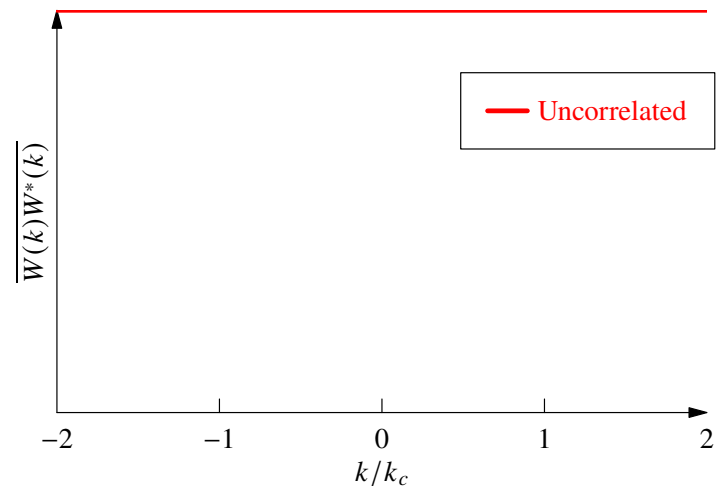
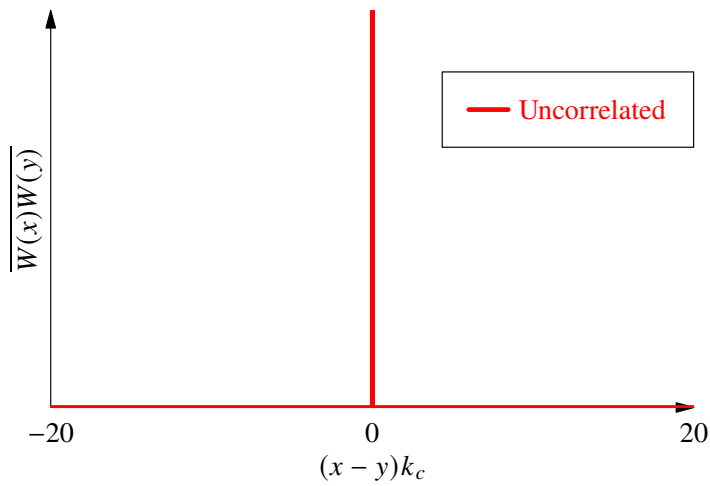
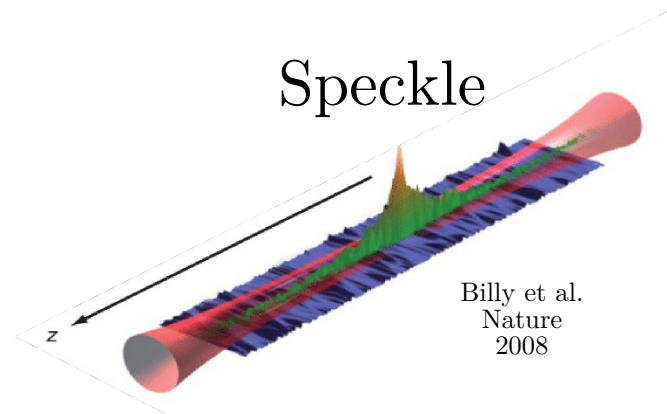
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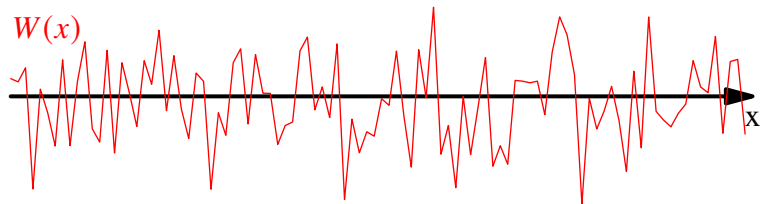
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Speckle



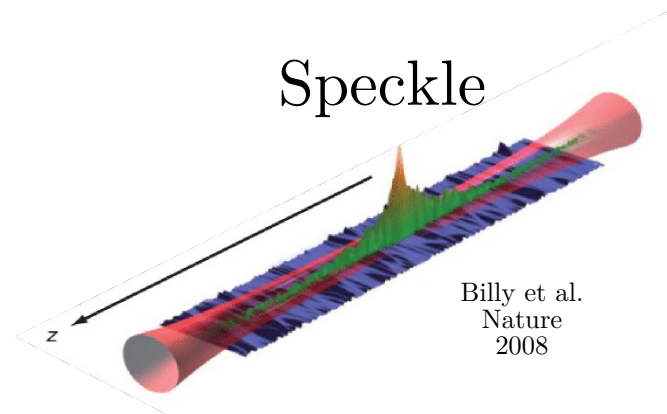
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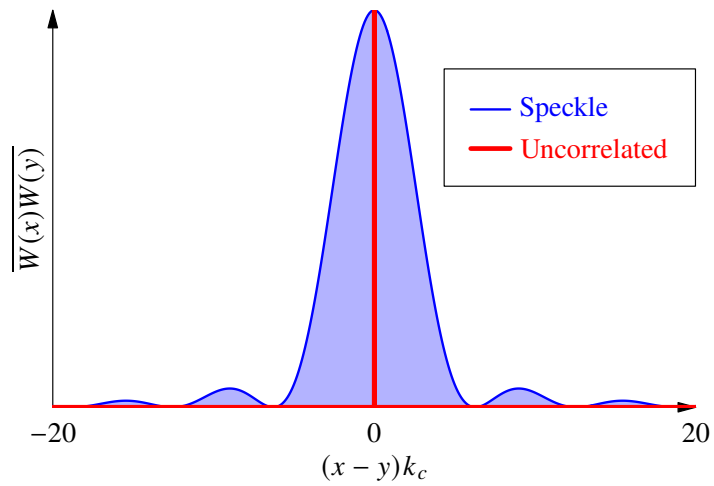


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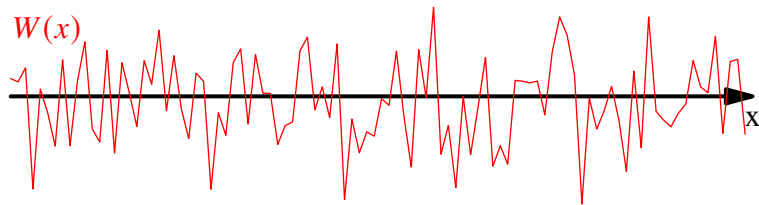


Billy et al.
Nature
2008



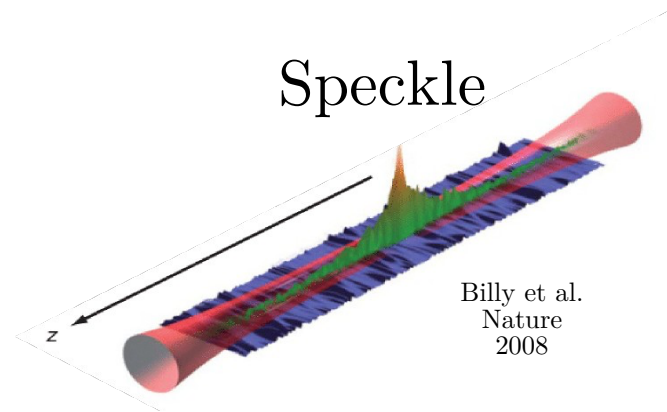
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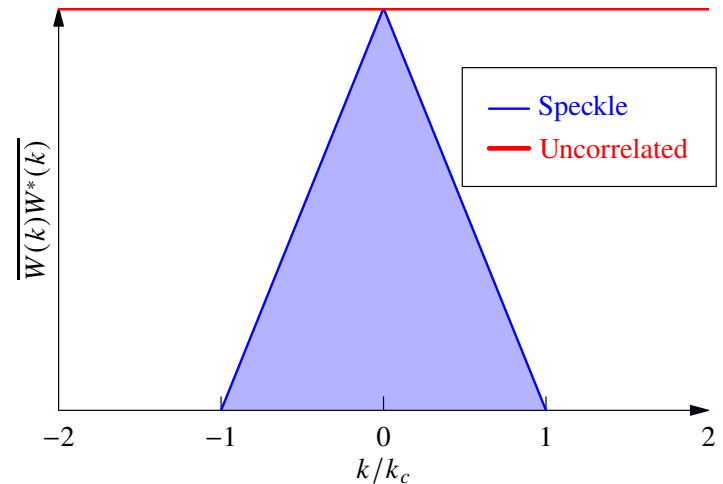
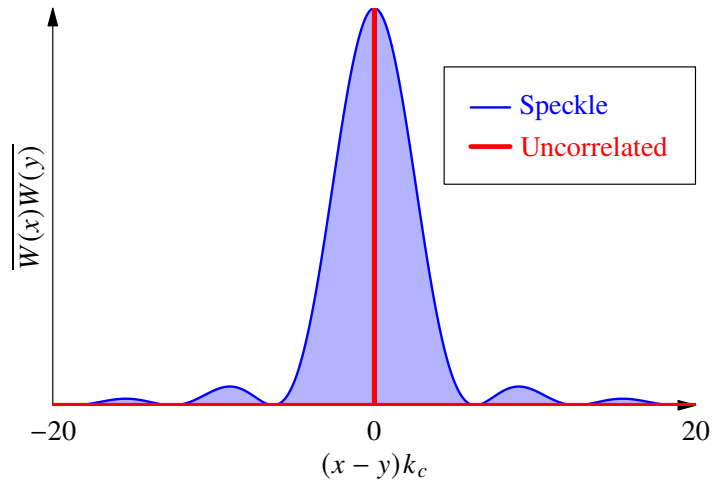


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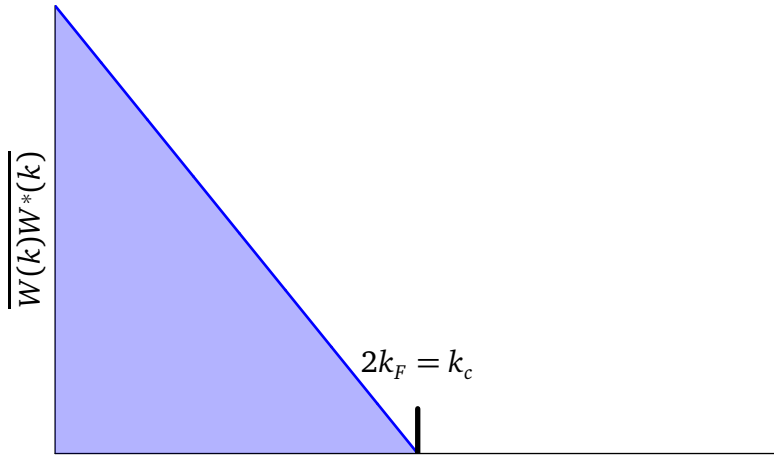
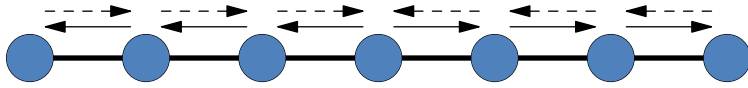
Speckle



$$\overline{W(k)W^*(k)} = W_0^2 \Omega (1 - |k|/k_c) \theta(k_c - |k|)$$



Our question



Backscattering barely allowed

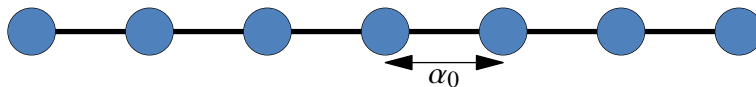
Are there other differences to uncorrelated disorder ?

What is the role of interactions ?

2 angles of attack

Bosonized Hamiltonian

Low-energy excitations are collective modes (in 1D)



$$H = \underbrace{\frac{1}{2\pi} \int dx \left(uK (\nabla\theta(x))^2 + \frac{u}{K} (\nabla\phi(x))^2 \right)}_{H_{\text{kin}} + H_{\text{int}}} + \underbrace{\frac{1}{2\pi\alpha_0} \int dx W(x) \cos(2k_F x - 2\phi(x))}_{H_{\text{dis}}}$$

Bosonic fields

$\phi(x)$ describes the density

$\theta(x)$ describes the phase

Interactions

K

Non-quadratic term
became quadratic

Disorder

Strength given by $W(x)$

Non-quadratic term

RG equations

Valid only for small
disorder strengths W_0

Integrating over the
degrees of freedom $\alpha = \alpha_0 e^l$

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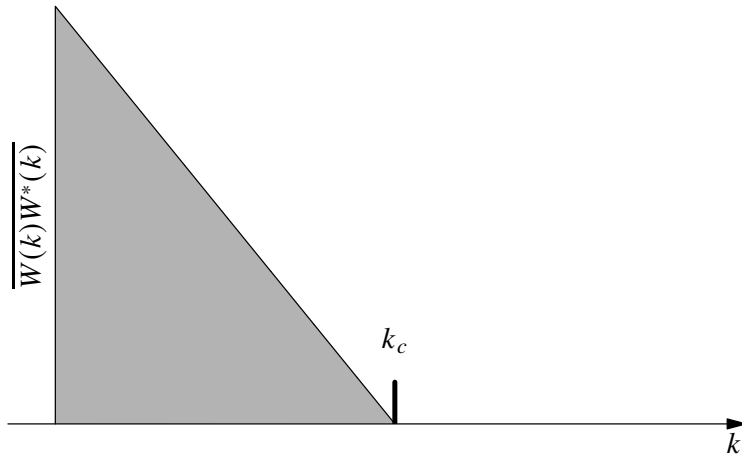
$$\frac{\partial K}{\partial l} \propto -\frac{K^2 W_0^2}{4} \sum_k (1 - |k|/k_c) \theta(k_c - |k|) [J_0((k + 2k_F)\alpha) + J_0((k - 2k_F)\alpha)]$$
$$\frac{\partial W_0^2}{\partial l} = (4 - 2K)W_0^2$$

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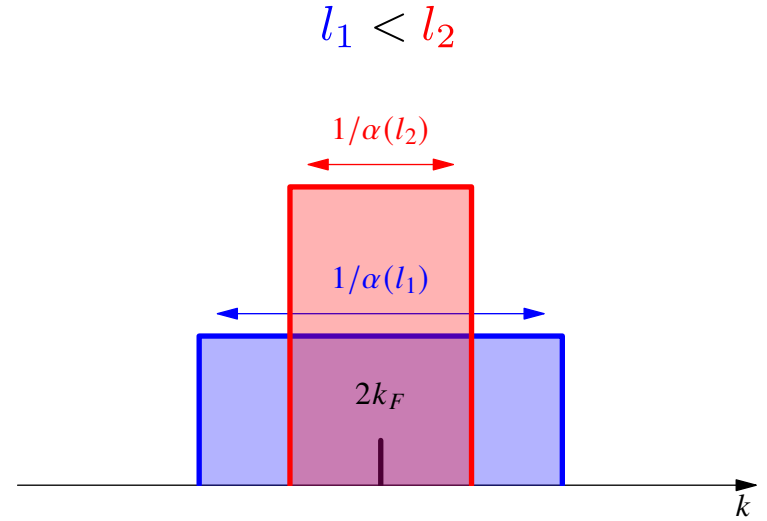
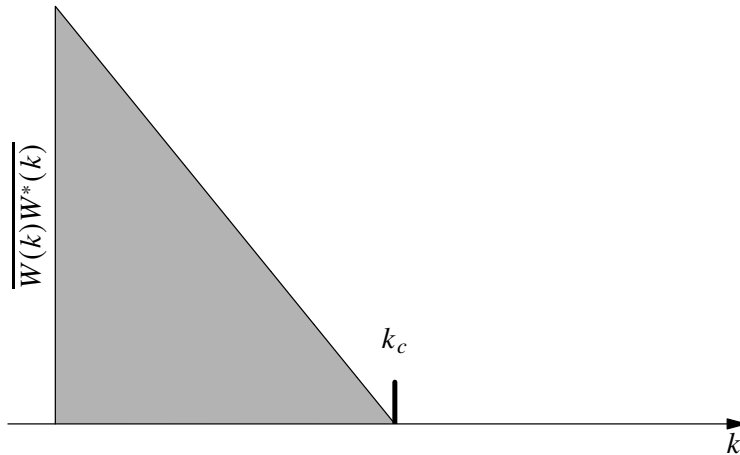
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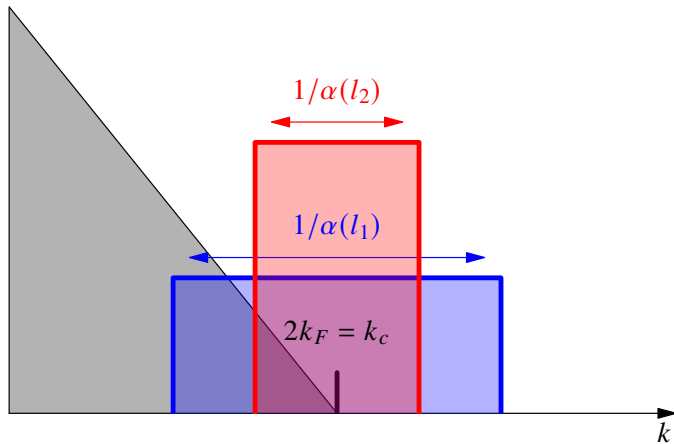
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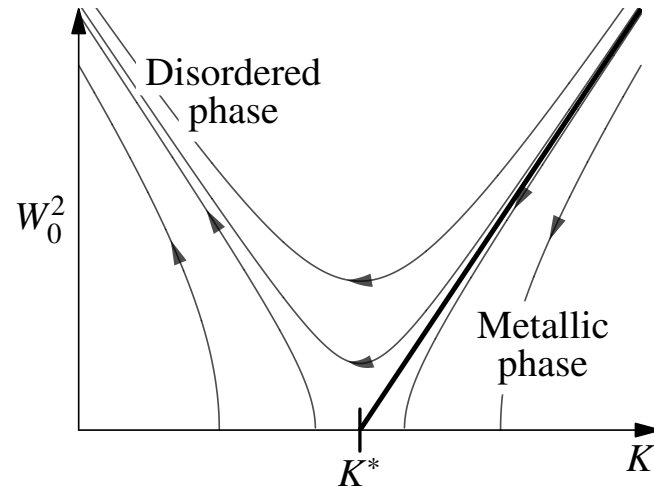
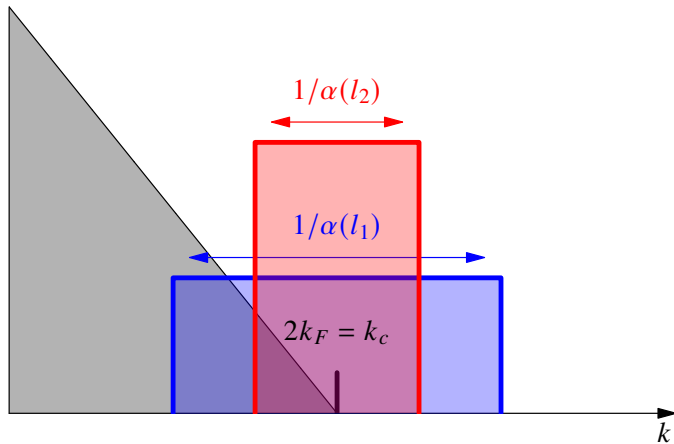
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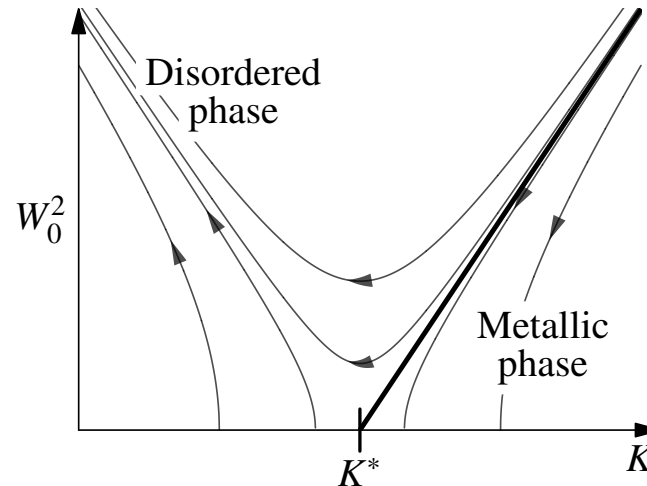
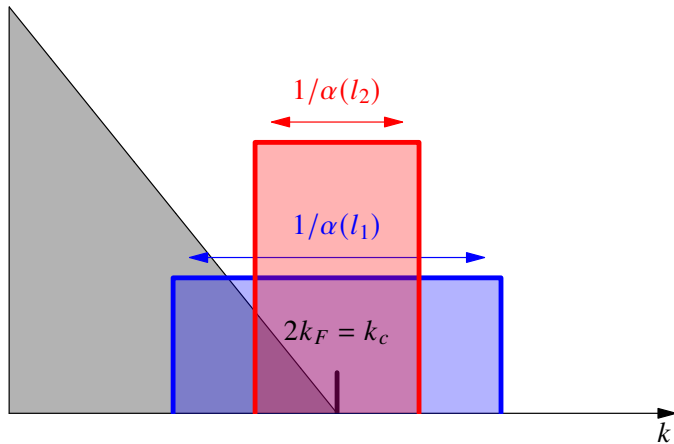
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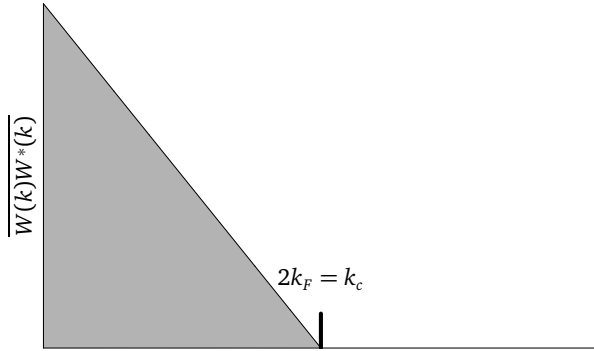


~~$K^* = 3/2$~~

$K^* = 1$

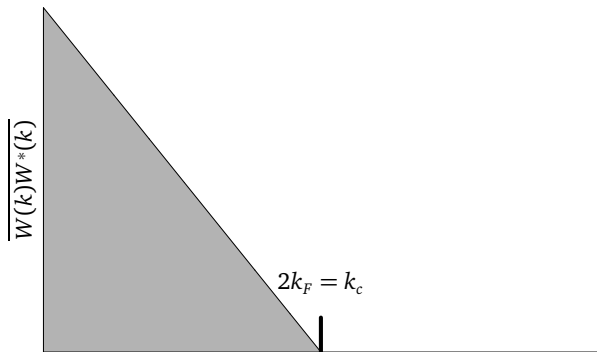
Shift of transition

K^* dependence of the correlations-shape

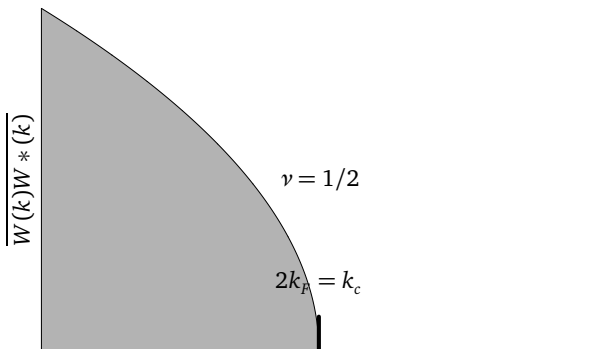
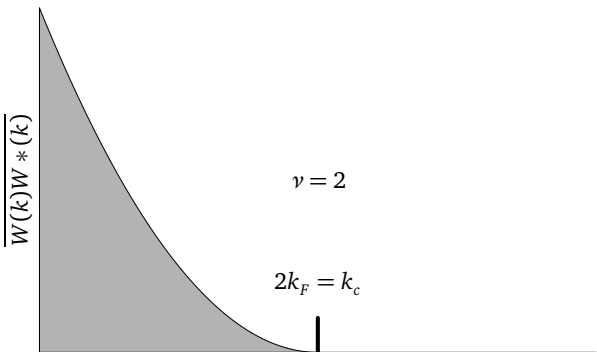


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K^* dependence of the correlations-shape

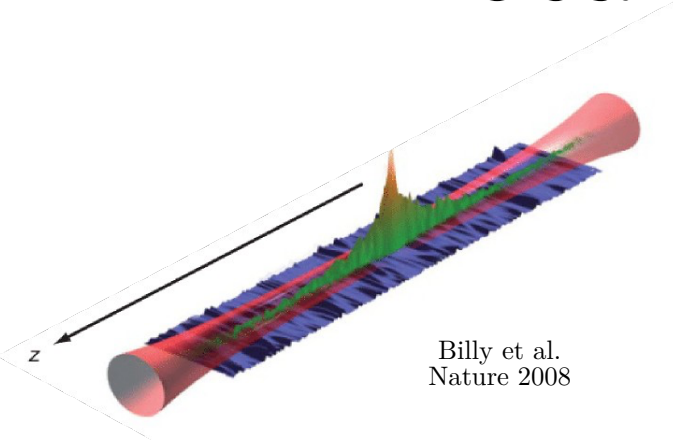


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$$K^* = \frac{3 - \nu}{2}$$

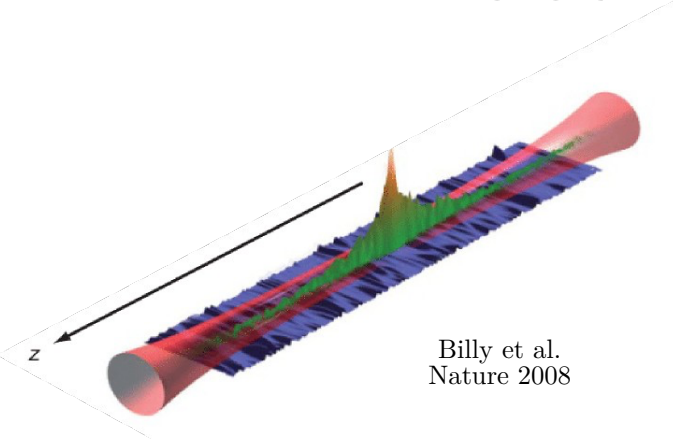
Localization length scaling



Billy et al.
Nature 2008

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Localization length scaling

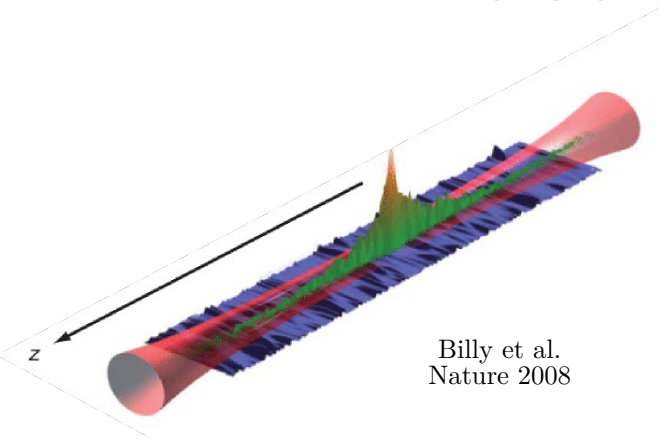


Billy et al.
Nature 2008

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Localization length scaling



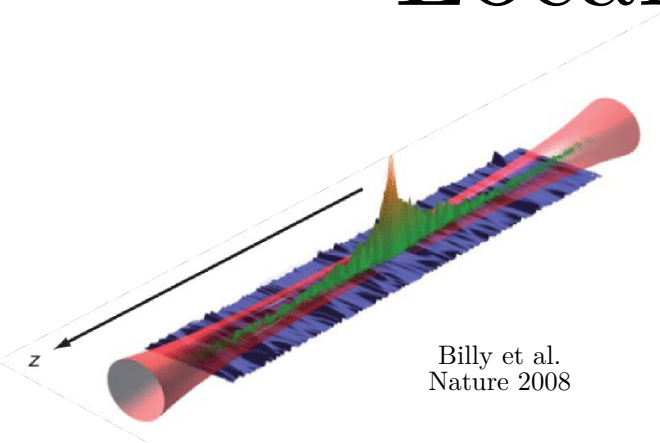
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Look at the
scaling of
 $\xi(W_0^2)$

Localization length scaling

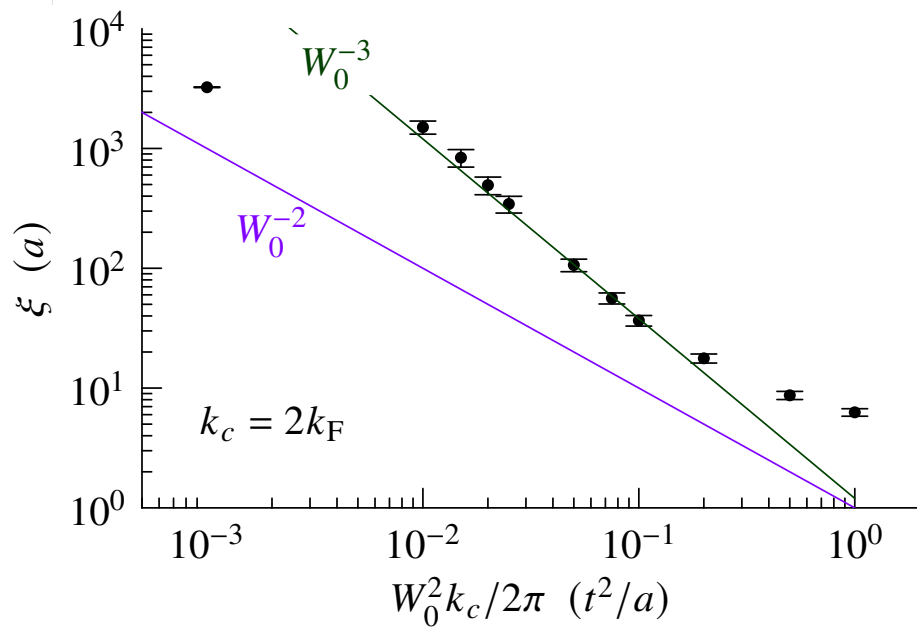


Billy et al.
Nature 2008

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Conclusions and outlook

- Vanishing backward processes lead to deviations from properties of uncorrelated disorder
 - Transition from a metallic phase to a disorder dominated phase is shifted to the non-interacting point
 - Deviation is also observed in the scaling of the localization length
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Thank you for your attention !

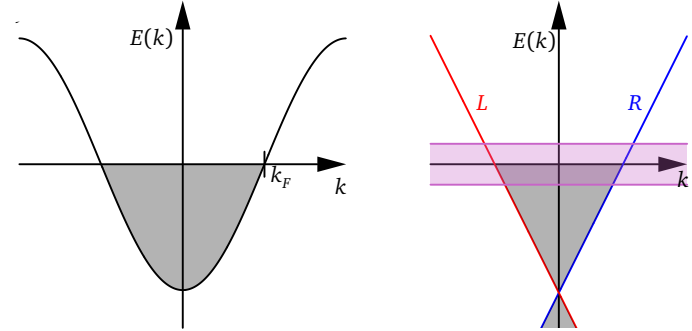
Hamiltonian

$$H = H_{\text{kin}} + H_{\text{int}} + H_{\text{dis}}$$

$$H_{\text{kin}} = \sum_r \sum_k v_F(\varepsilon_r k - k_F) c_{r,k}^\dagger c_{r,k}$$

$$H_{\text{int}} = \frac{g}{2\Omega} \sum_r \sum_{k,k',q} c_{r,k+q}^\dagger c_{-r,k'-q}^\dagger c_{-r,k'} c_{r,k}$$

$$H_{\text{dis}} = \frac{1}{\Omega} \sum_{k,q \sim 0} \left[W_{q-2k_F} c_{R,k+q}^\dagger c_{L,k} + W_{q+2k_F} c_{L,k+q}^\dagger c_{R,k} \right]$$

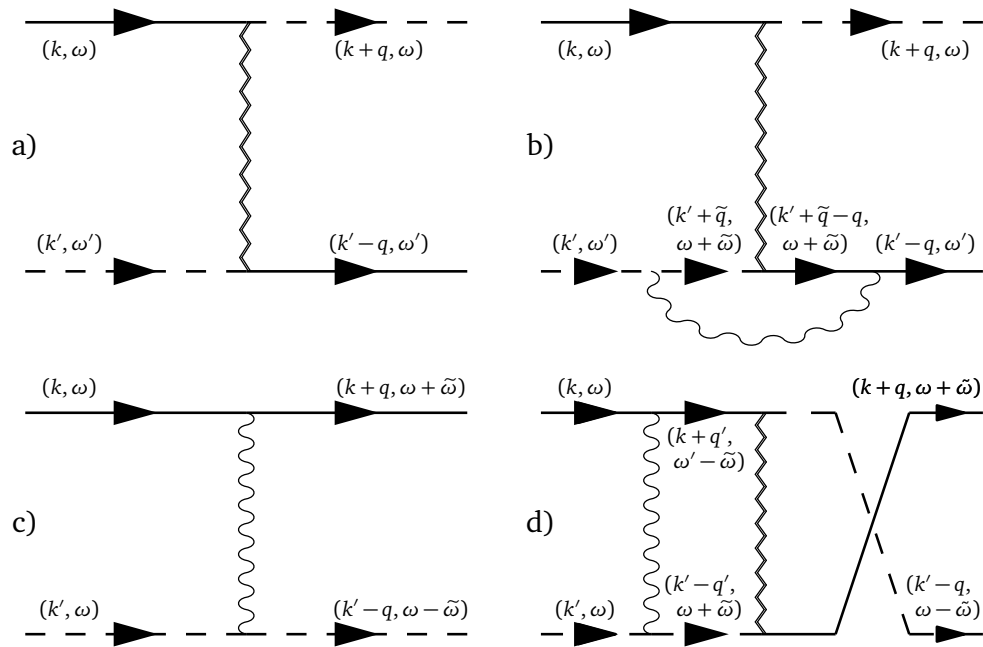


Linear kinetic part

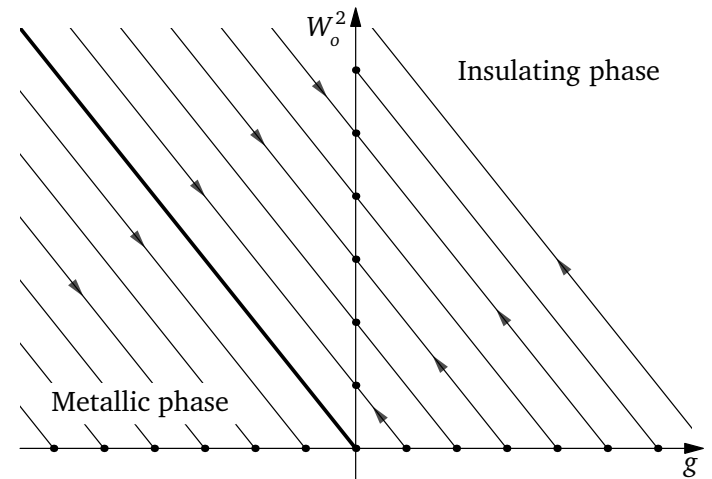
Interactions between left and right movers

Disorder processes change left into right movers

RG perturbative procedure



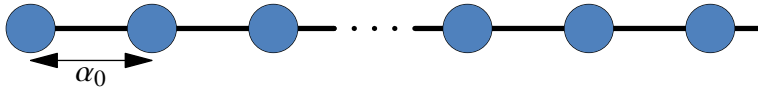
$$\begin{cases} \frac{\partial g}{\partial l} & \propto -gW_0^2 \\ \frac{\partial W_0^2}{\partial l} & \propto gW_0^2 \end{cases}$$



Renormalization Group (What is it?)

Idea :

- Complicated theory to investigate
- Interested only in macroscopic properties



- Integrate over some of the degrees of freedom
- Get a “new” theory (hopefully simpler)

- Repeat until you get a theory which can be treated

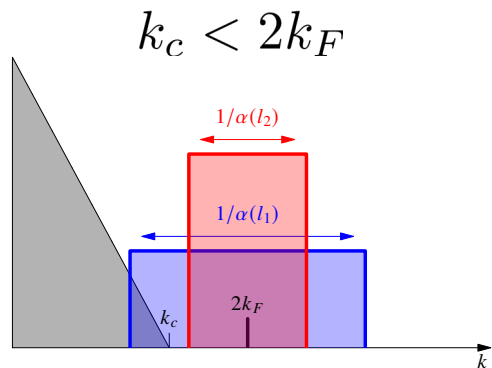
$$H(K, W_0)$$



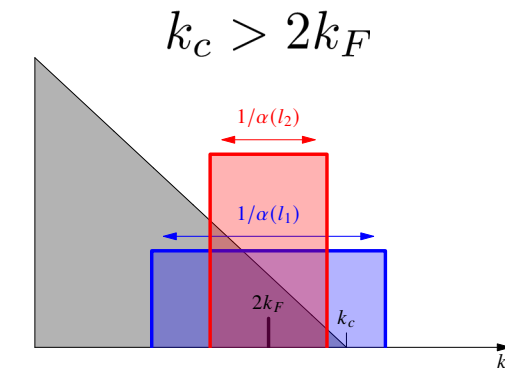
$$H^{\text{eff}}(K^{\text{eff}}, W_0^{\text{eff}}) = H(K^{\text{eff}}, W_0^{\text{eff}})$$

Looking for RG equations for the parameters of the theory K, W_0

RG equations : Flow of 3 cases



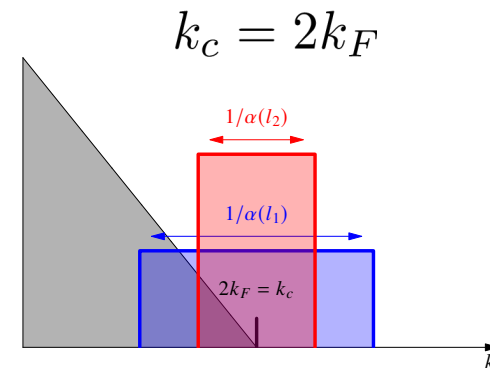
“Non-disordered”



$$\frac{\partial K}{\partial l} = -K^2 C e^{(3-2K)l}$$

Same behaviour as uncorrelated disorder

$$K^* = 3/2$$

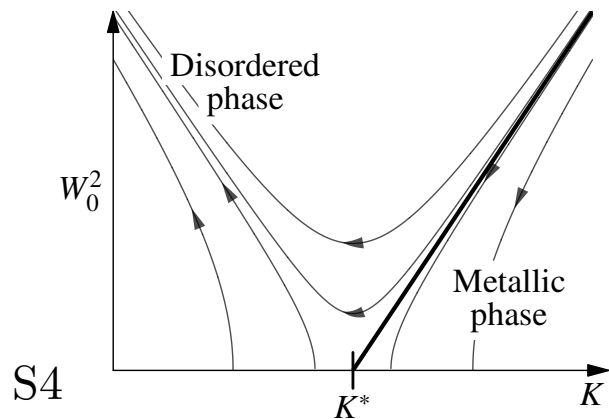


$$\frac{\partial K}{\partial l} = -K^2 C e^{(2-2K)l}$$

Different critical interaction

$$K^* = 1$$

Transition point shifted to non-interacting system



Localization length scaling

Compute ξ with IPR

$$\text{IPR}(E) = \int dx |\psi(x, E)|^4$$

